

# Trajectory Attractor for a System of Two Reaction-Diffusion Equations with Diffusion Coefficient $\delta(t) \rightarrow 0+$ as $t \rightarrow +\infty$ <sup>1</sup>

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The method of trajectory attractors (see [1–6]) makes it possible to effectively study the limit behavior of solutions to partial differential equations for which the theory of global attractors (see [7, 8]) is not directly applicable, for example, due to nonunique solvability of the corresponding Cauchy problem. In the present paper, we construct the trajectory attractor for a nonautonomous reaction-diffusion system where one equation has a time-dependent diffusion coefficient that tends to zero as  $t \rightarrow +\infty$ . In [9, 10], an analogous problem has been studied for autonomous reaction-diffusion systems having one time-independent diffusion coefficient that approaches zero as a parameter of the problem.

## 1. REACTION-DIFFUSION SYSTEM WITH A DIFFUSION COEFFICIENT $\delta(t)$ THAT APPROACHES ZERO AS $t \rightarrow +\infty$

In a bounded domain  $\Omega \subseteq \mathbb{R}^3$ , we consider the following nonautonomous reaction-diffusion system:

$$\partial_t u = \Delta u - f(u, v) + g_1(x), \quad (1)$$

$$\partial_t v = \delta(t) \Delta v - h(u, v) + g_2(x), \quad (2)$$

where  $u = u(x, t)$ ,  $v = v(x, t)$ ,  $x \in \Omega$ ,  $t \geq 0$ . The Laplace operator  $\Delta$  acts in the  $x$ -variable of the domain  $\Omega$ . At the boundary  $\partial\Omega$ , we set the homogeneous Dirichlet conditions

$$u|_{\partial\Omega} = 0, \quad v|_{\partial\Omega} = 0. \quad (3)$$

In Eq. (2), the diffusion coefficient  $\delta$  depends on time  $t$ .

We assume that  $\delta(\cdot) \in L_{\infty}^{\text{loc}}(\mathbb{R}_+)$ ,  $\delta(t) \geq 0$  for  $t \geq 0$ , and

$$\int_t^{t+1} \delta(s) ds \rightarrow 0 \quad \text{as } t \rightarrow +\infty. \quad (4)$$

The nonlinear functions  $f, h: \mathbb{R}^2 \rightarrow \mathbb{R}$  are continuous and satisfy the following inequalities:

$$\begin{aligned} \sigma_1(|u|^{p_1} + |v|^{p_2}) - c &\leq f(u, v)u + h(u, v)v \\ &\leq C(|u|^{p_1} + |v|^{p_2} + 1), \end{aligned} \quad (5)$$

$$\begin{aligned} |f(u, v)|^{p_1/(p_1-1)} + |h(u, v)|^{p_2/(p_2-1)} \\ \leq C(|u|^{p_1} + |v|^{p_2} + 1), \quad \forall u, v \in \mathbb{R} \end{aligned} \quad (6)$$

for some positive constants  $\sigma_1, c, C$ , and  $p_1, p_2 \geq 2$ . We also assume that  $h \in C^1(\mathbb{R}^2)$ ,  $h(0, 0) = 0$ , and the following inequalities hold:

$$0 < \sigma_2 \leq \frac{\partial h}{\partial v}(u, v), \quad \left| \frac{\partial h}{\partial u}(u, v) \right| \leq D, \quad \forall u, v \in \mathbb{R}. \quad (7)$$

The positive quantities  $\sigma_1$  and  $\sigma_2$  in (5) and (7) reflect the dissipation of the system (1)–(2). We set  $\sigma := \min\{\sigma_1, \sigma_2\}$ .

Note that we do not assume that the nonlinear function  $f$  in Eq. (1) is differentiable and so the Cauchy problem for system (1)–(3) can have more than one solution in the corresponding function space.

The functions  $g_1(x)$  and  $g_2(x)$  in Eqs. (1) and (2) belong to the spaces

$$g_1 \in L_2(\Omega), \quad g_2 \in H_0^1(\Omega).$$

An example of system (1), (2) is a nonautonomous analog of the FitzHugh–Nagumo system:

$$\partial_t u = \Delta u - u(u - \beta)(u - 1) - v,$$

$$\partial_t v = \delta(t) \Delta v + \alpha u - \gamma v,$$

where  $f(u, v) = u(u - \beta)(u - 1) + v$ ,  $h(u, v) = \gamma v - \alpha u$ , and  $\alpha, \beta, \gamma \geq 0$ . For this example,  $\sigma = \min\{1, \gamma\}$ ,  $p_1 = 4$ ,  $p_2 = 2$ , and  $g_1 \equiv 0$ ,  $g_2 \equiv 0$  (see [7, 11, 12]).

For simplicity of notations, we set  $H := L_2(\Omega)$ ,  $V := H_0^1(\Omega)$  and denote the norms in these spaces:  $\|u\| := \|u\|_H$ ,  $\|u\|_1 := \|u\|_V$ .

For arbitrary functions  $u(\cdot) \in L_{p_1}(0, M; L_{p_1}(\Omega))$  and  $v(\cdot) \in L_{p_2}(0, M; L_{p_2}(\Omega))$ , it follows from (6) that  $f(u, v) \in L_{q_1}(0, M; L_{q_1}(\Omega))$  and  $h(u, v) \in L_{q_2}(0, M; L_{q_2}(\Omega))$ , where  $q_i = \frac{p_i}{p_i - 1}$ ,  $i = 1, 2$ , and

<sup>1</sup> The article was translated by the authors.

$$\begin{aligned} & \|f(u, v)\|_{L_{q_1}(0, M; L_{q_1})}^{q_1} + \|h(u, v)\|_{L_{q_2}(0, M; L_{q_2})}^{q_2} \\ & \leq C_1(\|u\|_{L_{p_1}(0, M; L_{p_1})}^{p_1} + \|v\|_{L_{p_2}(0, M; L_{p_2})}^{p_2} + 1). \end{aligned}$$

**Definition 1.** A pair of functions  $(u(x, t), v(x, t))$ ,  $x \in \Omega, t \geq 0$ , is called a *weak solution* to system (1)–(3) if for every  $M > 0$ ,

$$u \in L_{p_1}(0, M; L_{p_1}(\Omega)) \cap L_2(0, M; V),$$

$$v \in L_{p_2}(0, M; L_{p_2}(\Omega)) \cap L_2(0, M; V),$$

and the functions  $u(x, t)$  and  $v(x, t)$  satisfy Eqs. (1) and (2) in the sense of distributions with values in the spaces  $H^{-r_1}(\Omega)$  and  $H^{-r_2}(\Omega)$ , respectively, where  $r_i =$

$\max\left\{1, 3\left(\frac{1}{2} - \frac{1}{p_i}\right)\right\}$ ,  $i = 1, 2$  (see [1, 8, 13]). The exponents  $r_i$  are chosen such that Eqs. (1), (2) and the Sobolev embedding theorem implies

$$\partial_t u(\cdot) \in L_{q_1}(0, M; H^{-r_1}(\Omega)),$$

$$\partial_t v(\cdot) \in L_{q_2}(0, M; H^{-r_2}(\Omega)).$$

If  $(u(\cdot), v(\cdot))$  is a weak solution to (1), (2), then it follows from these equations that

$$u(\cdot) \in L_\infty(0, M; H), \quad v(\cdot) \in L_\infty(0, M; H).$$

In addition, using the Lions–Magenes lemma (see [14]), we obtain

$$u(\cdot) \in C_w([0, M]; H),$$

$$v(\cdot) \in C_w([0, M]; H), \quad \forall M > 0.$$

Consequently, for every  $t \geq 0$ , the values of functions  $u(\cdot, t)$  and  $v(\cdot, t)$  are well defined as elements of the space  $H$ , and, in particular, the following initial conditions make sense:

$$u|_{t=0} = u_0 \in H, \quad v|_{t=0} = v_0 \in H. \quad (8)$$

In [1], we proved that for any weak solution  $(u(\cdot), v(\cdot))$  to system (1), (2), the functions  $u(\cdot)$  and  $v(\cdot)$  belong to  $C(\mathbb{R}_+; H)$ , while the function  $\|u(t)\|^2 + \|v(t)\|^2$  is absolutely continuous for  $t \geq 0$  and satisfies the energy identity

$$\begin{aligned} & \frac{1}{2} \frac{d}{dt} \{ \|u(t)\|^2 + \|v(t)\|^2 \} + \|\nabla u(t)\|^2 + \delta(t) \|\nabla v(t)\|^2 \\ & + \int_{\Omega} f(u(x, t), v(x, t)) u(x, t) dx \\ & + \int_{\Omega} h(u(x, t), v(x, t)) v(x, t) dx \end{aligned}$$

$$= \int_{\Omega} \{ g_1(x) u(x, t) + g_2(x) v(x, t) \} dx.$$

Formally, to obtain this identity, we multiply Eq. (1) by  $u(x, t)$  and Eq. (2) by  $v(x, t)$ . Then we integrate over  $x \in \Omega$  and sum the results.

**Proposition 1.** For any weak solution  $(u(\cdot), v(\cdot))$  to the system (1), (2) with initial data (8), the following inequalities hold

$$\begin{aligned} & \|u(t)\|^2 + \|v(t)\|^2 + 2 \int_0^t (\|u(s)\|_1^2 + \delta(s) \|v(s)\|_1^2) e^{-\sigma(t-s)} ds \\ & \leq (\|u_0\|^2 + \|v_0\|^2) e^{-\sigma t} + R_1^2, \end{aligned} \quad (9)$$

$$\begin{aligned} & 2 \int_0^{t+1} (\|u(s)\|_1^2 + \delta(s) \|v(s)\|_1^2) ds \\ & + \sigma \int_0^{t+1} (\|u(s)\|_{L_{p_1}}^{p_1} + \|v(s)\|_{L_{p_2}}^{p_2}) ds \\ & \leq (\|u_0\|^2 + \|v_0\|^2) e^{-\sigma t} + R_2^2, \quad \forall t \geq 0; \end{aligned} \quad (10)$$

the quantities  $R_1$  and  $R_2$  depend on  $\sigma$ ,  $\|g_1\|$ , and  $\|g_2\|$ .

We now assume that

$$v|_{t=0} = v_0 \in V. \quad (11)$$

Under this assumption, the existence of a weak solution to the problem (1)–(3), (8) is established using the Galerkin method.

**Proposition 2.** Under assumptions (11), the problem (1)–(3), (8) has a weak solution  $(u(\cdot), v(\cdot))$  such that  $v(\cdot) \in L_\infty(\mathbb{R}_+; V)$ , and the following inequality holds:

$$\begin{aligned} & \|v(t)\|_1^2 \leq \|v_0\|_1^2 e^{-\sigma t} + C_2(\|u_0\|^2 + \|v_0\|^2) e^{-\sigma t} + R^2, \\ & \forall t \geq 0, \end{aligned} \quad (12)$$

where  $R = R(\sigma, D, \|g_2\|, R_1)$ ,  $C_2 = C_2(\sigma, D)$ .

The Lions–Magenes lemma implies that  $v(\cdot) \in C_w(\mathbb{R}_+; V)$ ; i.e., the values  $v(\cdot, t) \in V$  are well-defined for all  $t \geq 0$ .

We now define a linear space  $\mathcal{F}_+^{\text{loc}}$ . By definition,

$$\mathcal{F}_+^{\text{loc}} = \left\{ \begin{aligned} & (y(x, t), z(x, t)), \quad x \in \Omega, t \geq 0 \\ & y \in L_\infty^{\text{loc}}(\mathbb{R}_+; H) \cap L_2^{\text{loc}}(\mathbb{R}_+; V) \cap L_{p_1}^{\text{loc}}(\mathbb{R}_+; L_{p_1}(\Omega)), \\ & z \in L_\infty^{\text{loc}}(\mathbb{R}_+; V) \cap L_{p_2}^{\text{loc}}(\mathbb{R}_+; L_{p_2}(\Omega)), \\ & \partial_t y \in L_{q_1}^{\text{loc}}(\mathbb{R}_+; H^{-r_1}(\Omega)), \quad \partial_t z \in L_{q_2}^{\text{loc}}(\mathbb{R}_+; H^{-r_2}(\Omega)) \end{aligned} \right\}. \quad (13)$$

We consider the linear subspace  $\mathcal{F}_+^b \subset \mathcal{F}_+^{\text{loc}}$  of functions  $(y, z)$  with a finite norm:

$$\begin{aligned} \|(y, z)\|_{\mathcal{F}_+^b} := & \|y\|_{L_\infty(\mathbb{R}_+; H)} + \|y\|_{L_2^b(\mathbb{R}_+; V)} + \|y\|_{L_{p_1}^b(\mathbb{R}_+; L_{p_1})} \\ & + \|\partial_t y\|_{L_{q_1}^b(\mathbb{R}_+; H^{-r_1})} + \|z\|_{L_\infty(\mathbb{R}_+; V)} + \|z\|_{L_{p_2}^b(\mathbb{R}_+; L_{p_2})} \\ & + \|\partial_t z\|_{L_{q_2}^b(\mathbb{R}_+; H^{-r_2})}. \end{aligned} \quad (14)$$

Recall that the norm in the space  $L_p^b(\mathbb{R}_+; X)$  is defined

$$\|y\|_{L_p^b(\mathbb{R}_+; X)} := \sup_{t \geq 0} \int_t^{t+1} \|y(s)\|_X^p ds.$$

The space  $\mathcal{F}_+^b$  with norm (14) is a Banach space.

We now define the space  $\mathcal{H}_+^\delta(N)$  of weak solutions (trajectories) to system (1), (2) with diffusion coefficient  $\delta$ , which depends on  $N > 0$ .

**Definition 2.** The space  $\mathcal{H}_+^\delta(N)$  consists of functions  $(u(\cdot), v(\cdot)) \in \mathcal{F}_+^{\text{loc}}$  such that (1) the pair  $(u(t), v(t))$ ,  $t \geq 0$  is a weak solution to (1), (2); (2) the function  $v(\cdot)$  satisfies the inequality

$$\|v(t)\|_1^2 \leq N e^{-\sigma t} + R^2, \quad \forall t \geq 0, \quad (15)$$

with  $\sigma$  and  $R$  from (12).

The space  $\mathcal{H}_+^\delta(N)$  is nonempty. Indeed, we take a weak solution to (1), (2) with initial data  $u_0 \in H$  and  $v_0 \in V$ , which was constructed in Proposition 2. If the norms  $\|u_0\|$ ,  $\|v_0\|_1$  satisfy inequality  $\|v_0\|_1^2 + C_2(\|u_0\|^2 + \|v_0\|_2^2) \leq N$ , then, due to (12), this weak solution belongs to  $\mathcal{H}_+^\delta(N)$ .

Consider the family of the translation operators  $T(\tau)$ ,  $\tau \geq 0$  acting on the space  $\mathcal{F}_+^{\text{loc}}$  by the formula

$$T(\tau)(y(t), z(t)) = (y(t + \tau), z(t + \tau)), \quad t \geq 0. \quad (16)$$

**Proposition 3.** The space of trajectories  $\mathcal{H}_+^\delta(N)$  lies in  $\mathcal{F}_+^b$ , and the following inequality holds:

$$\begin{aligned} \|T(\tau)(u, v)\|_{\mathcal{F}_+^b} & \leq C_3(\|u(0)\|^2 + N)e^{-\rho\tau} + R_3^2, \\ & \forall \tau \geq 0. \end{aligned} \quad (17)$$

## 2. THE LIMIT REACTION-DIFFUSION SYSTEM WITH ZERO DIFFUSION COEFFICIENT AND ITS TRAJECTORY ATTRACTOR

Consider the system that is “limit” with respect to (1), (2) as the diffusion coefficient  $\delta \equiv 0$ :

$$\partial_t u = \Delta u - f(u, v) + g_1(x), \quad (18)$$

$$\partial_t v = -h(u, v) + g_2(x). \quad (19)$$

This system is autonomous. At the boundary  $\partial\Omega$ , we set the conditions

$$u|_{\partial\Omega} = 0, \quad v|_{\partial\Omega} = 0. \quad (20)$$

By a *weak solution* to (18)–(20), we mean a pair of functions  $(u(x, t), v(x, t))$ ,

$$u(\cdot) \in L_{p_1}^{\text{loc}}(\mathbb{R}_+; L_{p_1}(\Omega)) \cap L_2^{\text{loc}}(\mathbb{R}_+; V),$$

$$v(\cdot) \in L_{p_2}^{\text{loc}}(\mathbb{R}_+; L_{p_2}(\Omega)) \cap L_\infty^{\text{loc}}(\mathbb{R}_+; V),$$

that satisfy Eqs. (18), (19) in the sense of the theory of distributions with values in the spaces  $H^{-r_1}(\Omega)$  and  $H^{-r_2}(\Omega)$ , respectively, (the exponents  $r_1$  and  $r_2$  have been defined in 1). For  $t = 0$ , we set the initial conditions

$$u|_{t=0} = u_0 \in H, \quad v|_{t=0} = v_0 \in V. \quad (21)$$

Propositions 1 and 2 are applicable to the system (18), (19) since it is a particular case of system (1), (2) for  $\delta \equiv 0$ . Therefore, the reduced inequalities (9), (10) are valid, where we remove terms containing the function  $\delta(s)$ . In addition, Cauchy problem (18)–(21) has a weak solution  $(u, v)$  that satisfies (15). From Eqs. (18), (19) we obtain

$$\partial_t u(\cdot) \in L_{q_1}(0, M; H^{-r_1}(\Omega)),$$

$$\partial_t v(\cdot) \in L_{q_2}(0, M; L_{q_2}(\Omega)).$$

Note that  $L_{q_2}(\Omega) \subset H^{-r_2}(\Omega)$  and, therefore, the constructed weak solutions belong to the space  $\mathcal{F}_+^{\text{loc}}$  (see (13)).

Similarly to the space  $\mathcal{H}_+^\delta(N)$ , we define the space of trajectories  $\mathcal{H}_+^0(N)$  for system (18), (19), which consists of its weak solutions  $(u(\cdot), v(\cdot))$  such that the function  $v(\cdot)$  satisfies inequality (15). As it was noticed before, the space  $\mathcal{H}_+^0(N)$  is nonempty and Proposition 3 holds: namely,  $\mathcal{H}_+^0(N)$  belongs to  $\mathcal{F}_+^b$ , and inequality (17) is true.

We now construct the trajectory attractor for system (18), (19).

In the space  $\mathcal{F}_+^{\text{loc}}$  (see (13)), we consider the following topology, which we define in terms of the weak convergence of the corresponding sequences from  $\mathcal{F}_+^{\text{loc}}$ . By definition, the sequence of functions  $\{(y_m(\cdot), z_m(\cdot))\} \subset \mathcal{F}_+^{\text{loc}}$  converges as  $m \rightarrow \infty$  to  $(y(\cdot), z(\cdot)) \in \mathcal{F}_+^{\text{loc}}$  in the topology  $\Theta_+^{\text{loc}}$  if, for each  $M > 0$ , the following convergences hold as  $m \rightarrow \infty$ :

(1)  $y_m(\cdot) \rightharpoonup y(\cdot)$  weak- $*$  in  $L_\infty(0, M; H)$ , weakly in  $L_2(0, M; V)$ , and weakly in  $L_{p_1}(0, M; L_{p_1}(\Omega))$ ;

(2)  $z_m(\cdot) \rightarrow z(\cdot)$  weak-\* in  $L_\infty(0, M; V)$  and weakly in  $L_{p_2}(0, M; L_{p_2}(\cdot))$ ;

(3)  $\partial_t y_m(\cdot) \rightarrow \partial_t y(\cdot)$  weakly in  $L_{q_1}(0, M; H^{-r_1}(\cdot))$ ;

(4)  $\partial_t z_m(\cdot) \rightarrow \partial_t z(\cdot)$  weakly in  $L_{q_2}(0, M; H^{-r_2}(\cdot))$ .

Note that the space  $\mathcal{F}_+^{\text{loc}}$  equipped with topology  $\Theta_+^{\text{loc}}$  is a linear Hausdorff space. In addition, this space is a Frechét–Urysohn space with a countable topology base (see, for example, [1]).

**Remark 1.** Any ball of radius  $r$

$$\mathcal{B}_r = \{ \|y, z\|_{\mathcal{F}_+^b} \leq r \}$$

in the space  $\mathcal{F}_+^b$  is a compact set in the topology  $\Theta_+^{\text{loc}}$ .

The set  $\mathcal{B}_r$  itself equipped with topology  $\Theta_+^{\text{loc}}$  is a compact metric space (see, for example, [15]).

Consider the translation operators  $T(\tau)$ ,  $\tau \geq 0$ , acting on  $\mathcal{F}_+^{\text{loc}}$  by the formula (16). These operators form a semigroup  $\{T(\tau)\} := \{T(\tau), \tau \geq 0\}$  on  $\mathcal{F}_+^{\text{loc}}$  called the *translation semigroup*. Note that  $\{T(\tau)\}$  takes the space of trajectories  $\mathcal{K}_+^0(N)$  to itself, since the system (18), (19) is autonomous; that is,

$$T(\tau): \mathcal{K}_+^0(N) \rightarrow \mathcal{K}_+^0(N), \quad \forall \tau \geq 0.$$

To construct the trajectory attractor, the key step is the following result proved in [10].

**Proposition 4.** *The space  $\mathcal{K}_+^0(N)$  is closed in  $\Theta_+^{\text{loc}}$  for every  $N \geq 0$ .*

Recall that the set  $P \subseteq \mathcal{K}_+^0(N)$  is called *absorbing* for the semigroup  $\{T(\tau)\}$  if, for every set  $B \subset \mathcal{K}_+^0(N)$  bounded in  $\mathcal{F}_+^b$ , there is  $\tau_1 = \tau_1(B) \geq 0$  such that  $T(\tau)B \subseteq P$  for all  $\tau \geq \tau_1$ . The set  $P \subseteq \mathcal{K}_+^0(N)$  is called *attracting* for the semigroup  $\{T(\tau)\}$  in the topology  $\Theta_+^{\text{loc}}$  if every neighbourhood  $\mathcal{O}(P)$  of the set  $P$  in the topology  $\Theta_+^{\text{loc}}$  is an absorbing set.

We now define the trajectory attractor for the semigroup  $\{T(\tau)\}|_{\mathcal{K}_+^0(N)}$ .

**Definition 3.** The set  $\mathfrak{A} \subset \mathcal{K}_+^0(N)$  is called the *trajectory attractor* for the semigroup  $\{T(\tau)\}$  on  $\mathcal{K}_+^0(N)$  if  $\mathfrak{A}$  is bounded in the norm  $\mathcal{F}_+^b$ , compact in the topology  $\Theta_+^{\text{loc}}$ , strictly invariant with respect to  $\{T(\tau)\}$ ,  $T(\tau)\mathfrak{A} = \mathfrak{A}$ , for all  $\tau \geq 0$ , and  $\mathfrak{A}$  is an attracting set for  $\{T(\tau)\}$  on  $\mathcal{K}_+^0(N)$  in the topology  $\Theta_+^{\text{loc}}$ .

We briefly describe the construction of the trajectory attractor. It follows from inequality (17) that the set

$$P_0 = \left\{ (u, v) \in \mathcal{K}_+^0(N) \mid \|(u, v)\|_{\mathcal{F}_+^b} \leq 2R_3^2 \right\}$$

is an absorbing set for the semigroup  $\{T(\tau)\}$  on  $\mathcal{K}_+^0(N)$  and, in addition,  $T(\tau)P_0 \subseteq P_0$  for all  $\tau \geq 0$ . As it was noticed before, the set  $P_0$  with topology  $\Theta_+^{\text{loc}}$ , being a ball in  $\mathcal{F}_+^b$ , is a compact metric space. It is easy to verify that the semigroup  $\{T(\tau)\}$  is continuous in the topology  $\Theta_+^{\text{loc}}$ . Consequently, we have a continuous semigroup  $\{T(\tau)\}$  acting on a compact metric space  $P_0$ . Then the general theorem on the existence of a global attractor is applicable (see, e.g., [7, 8]). The global attractor  $\mathfrak{A} \subseteq P$  of the semigroup  $\{T(\tau)\}|_{P_0}$  is constructed by the formula

$$\mathfrak{A} := \bigcap_{\tau \geq 0} \left[ \bigcup_{\theta \geq \tau} T(\theta)P_0 \right]_{\Theta_+^{\text{loc}}}.$$

The set  $\mathfrak{A}$  is clearly bounded in  $\mathcal{F}_+^b$ , compact in  $\Theta_+^{\text{loc}}$ , strictly invariant, that is,  $T(\tau)\mathfrak{A} = \mathfrak{A}$  for all  $\tau \geq 0$ , and it follows easily that  $\mathfrak{A}$  attracts any bounded in  $\mathcal{F}_+^b$  set  $B \subseteq \mathcal{K}_+^0(N)$  in the topology  $\Theta_+^{\text{loc}}$ . Therefore,  $\mathfrak{A}$  is the trajectory attractor of  $\{T(\tau)\}$  on  $\mathcal{K}_+^0(N)$ .

**Proposition 5.** *The trajectory attractor  $\mathfrak{A} = \mathfrak{A}(N)$  constructed above in the space  $\mathcal{K}_+^0(N)$  is independent of  $N$ . In particular,  $\mathfrak{A} = \mathfrak{A}(0)$ ; that is,*

$$\sup \{ \|v(t)\|_1^2 \mid t \geq 0 \} \leq R^2, \quad \forall (u, v) \in \mathfrak{A}.$$

### 3. CONVERGENCE OF SOLUTIONS OF THE ORIGINAL REACTION-DIFFUSION SYSTEM (1), (2) TO THE TRAJECTORY ATTRACTOR $\mathfrak{A}$ OF THE LIMIT SYSTEM (18), (19).

We now formulate the main theorem of the paper.

**Theorem 1.** *For any  $N > 0$  and for every set of trajectories  $B = \{(u(t), v(t))\} \in \mathcal{K}_+^0(N)$  of the nonautonomous system (1), (2), which is bounded in the space  $\mathcal{F}_+^b$ , the family of translations  $T(\tau)B = \{(u(t+\tau), v(t+\tau))\}$  converges in the topology  $\Theta_+^{\text{loc}}$  as  $\tau \rightarrow +\infty$  to the trajectory attractor  $\mathfrak{A}$  of the limit autonomous system (18), (19):*

$$T(\tau)B \rightarrow \mathfrak{A}(\tau \rightarrow +\infty) \text{ in } \Theta_+^{\text{loc}}.$$

In conclusion, we summarize the main result of the paper: the nonautonomous reaction-diffusion system (1), (2) with the diffusion coefficient  $\delta(t)$  that vanishes as  $t \rightarrow +\infty$  (in the sense (4)) has a trajectory attractor that coincides with the trajectory attractor  $\mathfrak{A}$  of the corresponding limit autonomous reaction-diffusion system (18), (19) with zero diffusion coefficient,  $\delta \equiv 0$ .

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